

The Riemann Problem

Systems of Conservation Laws

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Under the kind supervision of Dr. Tong Li

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Introduction to the Riemann Problem - Scalar Case

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- For $u \in \mathbb{R}^1$

$$u_t + f(u)_x = 0, \quad (1)$$

$$u(x, 0) = u_0(x) = \begin{cases} u_l, & x < 0 \\ u_r, & x > 0 \end{cases}$$

where u_l and u_r are constants.

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- This is in the scalar case of the Riemann problem.
- f must be genuinely nonlinear so WLOG, let $f'' > 0$.

Weak Solutions of Conservation Laws

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- A bounded measurable function $u(x, t)$ is called a *weak solution* of this IVP for any $\phi \in C_0^1(\mathbb{R}^1 \times [0, \infty))$

$$\int \int_{t \geq 0} (u \phi_t + f(u) \phi_x) dx dt + \int_{t=0} u_0 \phi dx = 0 \quad (2)$$

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- ϕ is called a test function.
- ϕ has compact support in $\mathbb{R}^1 \times [0, \infty)$.

Solution to the Riemann Problem - Shock Wave

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- The shock wave solution is:

$$u(x, t) = \begin{cases} u_l, & x < st \\ u_r, & x > st, \end{cases} \quad (3)$$

where

$$s = \frac{f(u_l) - f(u_r)}{u_l - u_r} \quad (4)$$

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- This is known as the jump condition (Rankine-Hugoniot condition)

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- Since we assumed $f'' > 0$ for all u , $u_l < u_r \Rightarrow f'(u_l) < f'(u_r)$.

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- If $f'(u_l) > s > f'(u_r)$, then we say that the entropy conditions are satisfied.
- If not, $f'(u_l) \not> s \not> f'(u_r)$, we have a rarefaction wave.
- Entropy condition **guarantees** uniqueness of weak solution.

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- We want to solve the Riemann Problem with n equations.
- Consider $\mathbf{u} \in \mathbb{R}^n$.
- $\mathbf{u} = (u_1, u_2, \dots, u_n)$, $\mathbf{f}(\mathbf{u}) = (f_1(\mathbf{u}), f_2(\mathbf{u}), \dots, f_n(\mathbf{u}))$
- The system

$$\mathbf{u}_t + \mathbf{f}(\mathbf{u})_x = \mathbf{0}$$

with initial data

$$\mathbf{u}(\mathbf{x}, 0) = \mathbf{u}_0(\mathbf{x}) = \begin{cases} \mathbf{u}_l, & x < 0 \\ \mathbf{u}_r, & x > 0 \end{cases} \quad (6)$$

is the Riemann problem with $\mathbf{u}_l, \mathbf{u}_r$ constant vectors.

System: Shock Waves

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- By (4), if \mathbf{u} has a discontinuity across $x = st$, the jump conditions need to be satisfied:

$$s[\mathbf{u}] = [\mathbf{f}(\mathbf{u})],$$

where $[\mathbf{u}] = \mathbf{u}_l - \mathbf{u}_r$ and $\mathbf{f}(\mathbf{u}) = \mathbf{f}(\mathbf{u}_l) - \mathbf{f}(\mathbf{u}_r)$

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- For systems, the *entropy inequalities* are as follows:

$$\lambda_k(\mathbf{u}_r) < s < \lambda_{k+1}(\mathbf{u}_r)$$

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for $1 \leq k \leq n$.

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- Such a discontinuity is called a *k-shock wave*.

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- In general, these class of equations have the following form:

$$\begin{cases} v_t - u_x = 0 \\ u_t + p(v)_x = 0, \quad t > 0, x \in \mathbb{R} \end{cases} \quad (7)$$

where $p' < 0, p'' > 0$.

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- In general, these class of equations have the following form:

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where $p' < 0, p'' > 0$.

- Rewrite the p -system as follows:

$$\mathbf{u}_t + \mathbf{f}(\mathbf{u})_x = \mathbf{0}$$

where $\mathbf{u} = (v, u)$ and $\mathbf{f}(\mathbf{u}) = (-u, p(v))$

Example

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- $p' < 0 \Rightarrow$ we have real and distinct eigenvalues \Rightarrow **strict hyperbolic system.**

Example (Continued)

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- Now, we can write the Riemann problem for the following initial conditions:

$$\mathbf{u}(\mathbf{x}, 0) = \mathbf{u}_0(\mathbf{x}) = \begin{cases} \mathbf{u}_l = (v_l, u_l), & x < 0 \\ \mathbf{u}_r = (v_r, u_r), & x > 0 \end{cases} \quad (8)$$

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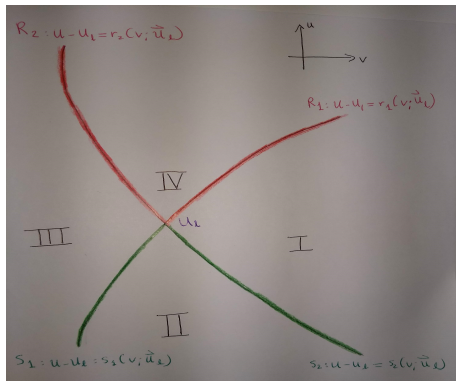
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- Entropy condition is satisfied.
- Given left state, what kind of right state can be connected to it?

Graphics



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- 2-shocks (front-shocks) satisfy

$$\lambda_1(\mathbf{u}_l) < s < \lambda_2(\mathbf{u}_l), \quad \lambda_2(\mathbf{u}_r) < s$$

- Hence, we have

$$-\sqrt{-p'(v_r)} < s < -\sqrt{-p'(v_l)}$$

and

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$$s = \frac{-(u - u_l)}{(v - v_l)} \text{ and } s = \frac{p(v) - p(v_l)}{u - u_l}$$

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$$s = \frac{-(u - u_l)}{(v - v_l)} \text{ and } s = \frac{p(v) - p(v_l)}{u - u_l}$$



$$-\frac{(u - u_l)}{(v - v_l)} = \frac{p(v) - p(v_l)}{u - u_l}$$

$$u - u_l = \pm \sqrt{(p(v_l) - p(v))(v - v_l)}$$

Continued

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- In order to form 1-shock, we need $-\sqrt{-p'(v)} < -\sqrt{p'(v_l)}$, which means $p'(v_l) > p'(v)$ and since $p'' > 0$, $v_l > v$

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- $S_1 : u - u_l = -\sqrt{(v - v_l)(p(v_l) - p(v))} \equiv s_1(v; \mathbf{u}_l)$, $v_l > v$

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- In order to form 1-shock, we need $-\sqrt{-p'(v)} < -\sqrt{p'(v_l)}$, which means $p'(v_l) > p'(v)$ and since $p'' > 0$, $v_l > v$
- $S_1 : u - u_l = -\sqrt{(v - v_l)(p(v_l) - p(v))} \equiv s_1(v; \mathbf{u}_l)$, $v_l > v$
- Similarly,
 $S_2 : u - u_l = -\sqrt{(v - v_l)(p(v_l) - p(v))} \equiv s_2(v; \mathbf{u}_l)$, $v_l < v$

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- A *rarefaction wave* is a continuous solution of the above system in the form $\mathbf{u} = U(x/t)$.
- We have 2 families of rarefaction waves, corresponding to either λ_1 or λ_2 .
- The k th family is genuinely nonlinear, $\nabla \lambda_k \cdot \mathbf{r}_k \neq 0$, where \mathbf{r}_k is the right eigenvalue.

Rarefaction - Continued

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- Let $\xi = x/t$

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- Hence, we have $\frac{du}{dv} = -\lambda_1(v, u) = \sqrt{-p'(v)}$
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- Similarly 2- rarefaction wave curve is given by
$$R_2 : u - u_l = - \int_{v_l}^v \sqrt{-p'(y)} dy \equiv r_2(v; \mathbf{u}_l), v_l > v$$

Conclude/ Describe Solution/ What It means

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- $S_1 : u - u_l = -\sqrt{(v - v_l)(p(v_l) - p(v))}, v_l > v$

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- $S_1 : u - u_I = -\sqrt{(v - v_I)(p(v_I) - p(v))}, v_I > v$
- $S_2 : u - u_I = -\sqrt{(v - v_I)(p(v_I) - p(v))}, v_I < v$
- $R_1 : u - u_I = \int_{v_I}^v \sqrt{-p'(y)} dy, v_I < v$

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Application: Isentropic gas dynamics model

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$$\begin{cases} v_t - u_x = 0 \\ u_t + \left(\frac{k}{v^\gamma}\right)_x = 0, \quad t > 0, x \in \mathbb{R} \end{cases}$$

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- v denotes the specific volume, i.e. $v = \rho^{-1}$, where ρ is the density, u denotes the velocity, and γ is the adiabatic gas constant.

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- Represents the conservation of mass and momentum
- v denotes the specific volume, i.e. $v = \rho^{-1}$, where ρ is the density, u denotes the velocity, and γ is the adiabatic gas constant.
- Note, in the p -system, if we choose $p(v) = kv^{-\gamma}$, we retrieve this isentropic gas dynamics equations.

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Definition

A *centered simple wave*, centered at (x_0, t_0) is a simple wave depending on $\frac{(x - x_0)}{(t - t_0)}$

Definition

The k th characteristic family is said to be *genuinely nonlinear* in a region $D \subset \mathbb{R}^n$ provided that $\nabla \lambda_k \cdot r_k \neq 0$ in D . If this is the case, normalize r_k by $\nabla \lambda_k \cdot r_k = 1$.

Solution of Riemann Problem for general hyperbolic systems

Theorem (Lax (1957))

Let $\mathbf{u}_l \in N \subset \mathbb{R}^n$.

Consider the system of n equations

$$\mathbf{u}_t + \mathbf{f}(\mathbf{u})_x = 0, x \in \mathbb{R}, t > 0,$$

where $\mathbf{u} = (u_1, \dots, u_n)$, $\mathbf{f}(\mathbf{u}) = (f_1(\mathbf{u}), f_2(\mathbf{u}), \dots, f_n(\mathbf{u}))$, the system is **hyperbolic.**, and each characteristic field is either genuinely nonlinear or linear degenerate in N .

Then, there is a neighborhood $\hat{N} \subset N$ of \mathbf{u}_l such that if $\mathbf{u}_r \in \hat{N}$, the Riemann problem has precisely one solution, consisting of at most $(n + 1)$ constant states.

Application To Gas Dynamics

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$$v_t - u_x = 0$$

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$$\begin{pmatrix} 0 & -1 & 0 \\ p_v & 0 & p_s \\ 0 & 0 & 0 \end{pmatrix}$$

Application To Gas Dynamics

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$$v_t - u_x = 0$$

$$u_t + p_x = 0$$

$$(e + 1/2u^2)_t + (pu)_x = 0$$

-

$$\begin{pmatrix} 0 & -1 & 0 \\ p_v & 0 & p_s \\ 0 & 0 & 0 \end{pmatrix}$$

- The eigenvalues are $\lambda_1 = -\sqrt{-p_v}$, $\lambda_2 = 0$, $\lambda_3 = \sqrt{-p_v}$

Application To Gas Dynamics



$$v_t - u_x = 0$$

$$u_t + p_x = 0$$

$$(e + 1/2u^2)_t + (\rho u)_x = 0$$



$$\begin{pmatrix} 0 & -1 & 0 \\ p_v & 0 & p_s \\ 0 & 0 & 0 \end{pmatrix}$$

- The eigenvalues are $\lambda_1 = -\sqrt{-p_v}$, $\lambda_2 = 0$, $\lambda_3 = \sqrt{-p_v}$
- We have two *genuinely nonlinear* characteristic families, and one linearly degenerate. So, we only have 2 families of shock waves and rarefaction waves.

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References

1. Borovikov, V. On the problem of discontinuity decay for a system of two quasilinear equations. *Dokl. Akad. , SSSR*, **185** (1969), 250-252; English transl. in *Sov. Math., Dokl.*, **10** (1969), 321-323.
2. Cooper, Jeffrey. *Introduction to Partial Differential Equations with MATLAB*. New York. Springer Science. 1998.
3. Lax, P. Hyperbolic systems of conservation laws, II. *Comm. Pure Appl. Math.*, **10** (1957), 537-566.
4. Smoller, J. 1994. *Shock Waves and Reaction–Diffusion Equations*. New York. Springer-Verlag.
5. Smoller, J. On the solution of Riemann problem with general step data for an extended class of hyperbolic systems. *Mich. Math. J.*, **16** (1969), 201-210.

The End

- Thank You!
- Questions?