

Global BV solution to a system of balance laws from traffic flow

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Traffic System

- ▶ Constructing global solutions and finding zero relaxation limits of traffic flow
- ▶ Roadways, Vehicles, Drivers
- ▶ Microscopic Vs Macroscopic
- ▶ We will be focusing on a specific macroscopic model

History of Traffic Flow

- ▶ Lighthill-Whitham-Richards (LWR) model [1955, 1956]
- ▶ Payne-Whitham (PW) model [1971, 1974]
- ▶ Viscous models studied by Kerner-Konhauser, Kühne, Beckshulte, and Li [1984-1994, 2008]
- ▶ Aw-Rascle and Zhang's higher continuum (ARZ) models [2000, 2001]
- ▶ There are more references not mentioned above

Nonlinear Balance Laws

- ▶ Let $U \in \mathbb{R}^n$.
- ▶ $U = (u_1, u_2, \dots, u_n)$, $F(U) = (f_1(u), f_2(u), \dots, f_n(u))$
- ▶ Consider the general conservation form

$$U_t + F(U)_x + P(U) = 0 \quad (1)$$

with initial data

$$U(x, 0) = U_0(x), \quad (2)$$

where $x \in \mathbb{R}$, $t > 0$.

The Model

- ▶ To analyze the 2×2 traffic flow model:

$$\begin{aligned}\rho_t + (\rho v)_x &= 0, \\ v_t + \left(\frac{1}{2}v^2 + g(\rho)\right)_x + \frac{v - v_e(\rho)}{\tau} &= 0,\end{aligned}\tag{3}$$

with initial data

$$(\rho(x, 0), v(x, 0)) = (\rho_0(x), v_0(x))\tag{4}$$

where $x \in \mathbb{R}$, $t > 0$, $\tau > 0$.

- ▶ ρ - density, v - velocity, $v_e(\rho)$ - equilibrium velocity.
- ▶ $g(\rho)$ - anticipation factor and satisfies

$$g'(\rho) = \rho(v_e'(\rho)/\theta)^2,\tag{5}$$

where $g'(\rho) \geq 0$, $0 < \theta < 1$.

LWR Model

- ▶ The equilibrium flow is described by Lighthill-Whitham-Richards (LWR) model [8, 9]

$$\rho_t + (\rho v_e(\rho))_x = 0, \quad x \in \mathbb{R}, t > 0, \quad (6)$$

with initial data $\rho(x, 0) = \rho_0(x) > 0$.

- ▶ $q(\rho) = \rho v_e(\rho)$ is known as the fundamental diagram
- ▶ For our work, we let

$$v_e(\rho) = -a\rho + b, \quad (7)$$

where $a > 0, b > 0$.

- ▶ In our study, the equilibrium flux $q(\rho) = \rho(-a\rho + b)$ is a concave function of ρ .

Work Overview

- ▶ We showed in [1] the existence of a global BV solution for a system of balance laws arising in traffic flow in the framework of Dafermos [2]
- ▶ Computed entropy-entropy flux pair, Kawashima condition, sub-characteristic condition, and the partial dissipative inequality
- ▶ With these conditions we show the existence of BV solutions for the Cauchy problem

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First Transformation

- ▶ We want $U \equiv 0$ to be an equilibrium solution we need to do a change of variables $v = u + b$.
- ▶ Now we can rewrite (3) as follows

$$\begin{aligned} \rho_t + (\rho(u + b))_x &= 0 \\ u_t + \left(\frac{1}{2}(u + b)^2 + g(\rho)\right)_x + \frac{u + b - v_e(\rho)}{\tau} &= 0 \end{aligned} \quad (8)$$

- ▶ In terms of the general form, we have

$$\begin{aligned} U &= (\rho, v - b) = (\rho, u)^T \\ F(U) &= (\rho(u + b), \frac{1}{2}u^2 + ub + g(\rho))^T \\ P(U) &= (0, \frac{u + b - v_e(\rho)}{\tau})^T \end{aligned} \quad (9)$$

Preliminaries

- ▶ The Jacobian is

$$\begin{bmatrix} u + b & \rho \\ g'(\rho) & u + b \end{bmatrix} \quad (10)$$

- ▶ Using (5) and (7), the eigenvalues are

$$\lambda_{1,2} = u + b \mp \frac{a}{\theta} \rho \quad (11)$$

- ▶ The corresponding right eigenvectors are

$$r_{1,2} = \left(\mp \frac{\theta}{a}, 1 \right)^T. \quad (12)$$

- ▶ The system (8) is genuinely nonlinear since

$$\nabla \lambda_i \cdot r_i = \frac{q''(\rho)}{v_e'(\rho)} = 2 \neq 0, \quad i = 1, 2. \quad (13)$$

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Obstacles

In order to apply Dafermos' theory [2], we had to

- ▶ Search for a convex entropy-entropy flux pair
- ▶ Verify conditions
- ▶ Transform system (3) once again into equivalent form

Entropy-Entropy Flux Pair

- ▶ We need to find smooth entropy flux pair $(\eta, q)(U)$ where η is convex and has been normalized by $\eta(0)$, $D\eta(0) = 0$.
- ▶ This is important since admissible solutions U must satisfy the entropy inequality

$$\partial_t \eta(U(x, t)) + \partial_x q(U(x, t)) + D\eta(U(x, t))P(U(x, t)) \leq 0 \quad (14)$$

- ▶ We also want our system to be a **symmetrizable**, which means it needs to be endowed with nontrivial **companion balance laws**.
- ▶ So we also need to solve

$$\begin{aligned} DQ_1(U, X) &= B(U, X)^T DG_1(U, X) \\ DQ_2(U, x) &= B(U, X)^T DG_2(U, X), \end{aligned} \quad (15)$$

where $G_1 = U$, $G_2 = F(U)$, $DQ_i = \left[\frac{\partial Q_i}{\partial \rho}, \frac{\partial Q_i}{\partial u} \right]$, $i = 1, 2$.

- ▶ Solving (15), we then constructed an explicit solution of a convex entropy-entropy flux pair

$$\eta(\rho, u) = Q_1(\rho, u) = (u - s\rho)^2 + \Gamma(u + s\rho)^2, \quad (16)$$

$$\begin{aligned} q(\rho, u) = Q_2(\rho, u) = & ((u - s\rho)^2 + \Gamma(u + s\rho)^2)(u + b) \\ & + (1 + \Gamma)(u(s\rho)^2) \\ & + 2(\Gamma - 1)\frac{(s\rho)^3}{3} - \frac{1 + \Gamma}{3}u^3, \end{aligned} \quad (17)$$

where $s = \frac{a}{\theta}$, $\Gamma = \frac{1 + \theta}{1 - \theta} > 1$.

- ▶ With this entropy-entropy flux pair, the convexity conditions are satisfied

Partial Dissipative Inequality

- ▶ We assume that P is dissipative semidefinite relative to η , i.e.

$$D\eta(U) \cdot P(U) \geq \alpha |P(U)|^2, \quad (18)$$

with $\alpha > 0$.

- ▶ For our system (8), we needed to find a condition such that

$$\begin{bmatrix} \frac{\partial \eta}{\partial \rho} & \frac{\partial \eta}{\partial u} \end{bmatrix} \cdot \begin{bmatrix} 0 \\ \frac{u + b - v_e(\rho)}{\tau} \end{bmatrix} \geq \alpha \left(\frac{u + b - v_e(\rho)}{\tau} \right)^2 \quad (19)$$

- ▶ After simplification, we require

$$0 < \alpha \leq \tau(2\Gamma + 1), \quad (20)$$

where $\Gamma > 1$.

Kawashima Condition

- ▶ The Kawashima condition is given by

$$DP(0)r_i(0) \neq 0, \quad i = 1, 2. \quad (21)$$

- ▶ For our system, we have

$$DP(0)r_i(0) = \begin{bmatrix} 0 \\ \frac{\mp\theta + 1}{\tau} \end{bmatrix} \neq \begin{bmatrix} 0 \\ 0 \end{bmatrix} \quad (22)$$

since $0 < \theta < 1$.

Sub-characteristic condition

- ▶ The sub-characteristic is satisfied when

$$\lambda_1 < \lambda_* < \lambda_2. \quad (23)$$

- ▶ For $v = v_e(\rho)$,

$$\lambda_*(\rho) = -2a\rho + b.$$

- ▶ The sub-characteristic condition is satisfied for (8) since we have

$$v_e(\rho) - \frac{a}{\theta}\rho < -2a\rho + b < v_e(\rho) + \frac{a}{\theta}\rho \quad (24)$$

for $0 < \theta < 1$.

Equivalent Form

- ▶ In order to apply Dafermos' theory, we needed to convert (8) into an equivalent form

$$\begin{aligned}\partial_t V + \partial_x G(V, W) + X(V, W) &= 0 \\ \partial_t W + \partial_x H(V, W) + CW + Y(V, W) &= 0,\end{aligned}\tag{25}$$

where $x \in \mathbb{R}$, $t > 0$, and $\eta_{WW}C(0, 0) > 0$.

- ▶ We followed Dafermos [2] and found the following change of variables

$$Z = (V, W) = (\rho, a\rho + u),\tag{26}$$

which transforms (8) to

$$\begin{aligned}V_t + [V(W - aV + b)]_x &= 0 \\ W_t + \left[\frac{1}{2}(W^2 - a^2V^2) + bW + g(V)\right]_x + \frac{1}{\tau}W &= 0\end{aligned}\tag{27}$$

with initial conditions

$$Z_0 = (V_0, W_0) = (\rho_0, a\rho_0 + u_0).\tag{28}$$

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Theorem (Admissible BV solution to the Cauchy problem)

Consider the Cauchy problem (27), (28) with genuinely nonlinear characteristic fields (13), endowed with a convex entropy η (16) (17). The source is dissipative semidefinite (18) relative to the entropy η , and the Kawashima condition (21) holds. Then there are positive constants $\delta_1, \sigma_0, c_0, c_1, \nu, b$ so that the Cauchy problem (27) (28) under initial data Z_0 with

$$\int_{-\infty}^{\infty} (1 + x^2) |Z_0(x)|^2 dx = \sigma^2 < \sigma_0^2, \quad (29)$$

$$TV_{(-\infty, \infty)} Z_0(\cdot) = \delta < \delta_1, \quad (30)$$

$$\int_{-\infty}^{\infty} V_0(x) dx = 0, \quad (31)$$

possesses an admissible BV solution Z on $(-\infty, \infty) \times [0, \infty)$ and

Theorem (Continued)

$$\int_{-\infty}^{\infty} |Z(x, t)| dx \leq b\sigma, \quad 0 \leq t < \infty, \quad (32)$$

$$TV_{(-\infty, \infty)} Z(\cdot, t) \leq c_0\sigma + c_1\delta e^{-\nu t}, \quad 0 \leq t < \infty, \quad (33)$$

$$\int_{-\infty}^{\infty} |Z(x, t)| dx \rightarrow 0, \quad \text{as } t \rightarrow \infty, \quad (34)$$

$$TV_{(-\infty, \infty)} Z(\cdot, t) \rightarrow 0, \quad \text{as } t \rightarrow \infty, \quad (35)$$

where $\delta > 0$.

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